

ON THE DARK MATTER CONTENT OF ULTRA-DIFFUSE GALAXIES

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ABSTRACT

I compare the dark matter content within stellar half-mass radius expected in a Λ CDM-based galaxy formation model with existing observational estimates for the observed dwarf satellites of the Milky Way and ultra-diffuse galaxies (UDGs). The model reproduces the main properties and scaling relations of dwarf galaxies, in particular their stellar mass-size relation. I show that the model also reproduces the relation between the dark matter mass within the half-mass radius, $M_{\text{dm}}(< r_{1/2})$, and stellar mass exhibited by the observed dwarf galaxies. The scatter in the $M_{\text{dm}}(< r_{1/2}) - M_{\star}$ relation is driven primarily by the broad range of sizes of galaxies of a given stellar mass. I also show the $M_{\text{dm}}(< r_{1/2})$ of UDGs are within the range expected in the model for their stellar mass, but they tend to lie above the median relation due to their large sizes. The upper limits on $M_{\text{dm}}(< r_{1/2})$ for the dark matter deficient UDGs are also consistent with the range of dark matter masses expected in the model. The most dark matter-deficient galaxies of a given size correspond to halos with the smallest concentrations and the largest ratios of M_{\star}/M_{200c} . Conversely, the most dark matter-dominated galaxies are hosted by the highest concentration halos with the smallest M_{\star}/M_{200c} ratios. The model indicates that the scatter between $M_{\text{dm}}(< r_{1/2})$ and M_{200c} is large, which renders inference of the virial mass from $M_{\text{dm}}(< r_{1/2})$ uncertain and dependent on specific assumptions about the halo mass profile. Results presented in this paper indicate that dark matter-deficient UDGs may represent a tail of the expected dark matter profiles, especially if the effect of feedback on these profiles is taken into account.

Subject headings: galaxies: structure, galaxies: dark matter, galaxies: halos

1. INTRODUCTION

The advent of deep imaging with CCD cameras resulted in discovery and investigations of the low surface brightness galaxies (e.g., Bothun et al. 1985, 1987; Impey et al. 1988; Schombert et al. 1992; Sprayberry et al. 1995; McGaugh et al. 1995; Dalcanton et al. 1997; Bothun et al. 1997; Impey & Bothun 1997; Driver et al. 2005; Geller et al. 2012), confirming early conclusions of Disney (1976) that strong selection effects bias galaxy studies against low surface brightness objects.

Studies of such galaxies have been re-energized in the last decade by the discovery of hundreds of low-surface brightness galaxies (re-branded as ultra-diffuse galaxies or UDGs) in the outskirts of the Coma galaxy cluster (van Dokkum et al. 2015a,b; Koda et al. 2015; Yagi et al. 2016a; Bautista et al. 2023) and other nearby clusters (Conselice et al. 2003; van der Burg et al. 2016, 2017; Román & Trujillo 2017; Venhola et al. 2017, 2022; Mancera Piña et al. 2018, 2019; Gannon et al. 2022a; Forbes et al. 2023; Zöllner et al. 2024; Marleau et al. 2024), although indications that such galaxies are present in the Virgo cluster date back to Sandage & Binggeli (1984) and Impey et al. (1988). Ultra-diffuse galaxies have also been found in the field (e.g., Leisman et al. 2017). These gas-rich field UDGs are likely to be evolutionarily linked to dwarf irregular galaxies (Nandi et al. 2023). UDGs in clusters also represent a tail of the general dwarf spheroidal population (Trujillo 2021; Nandi et al. 2023; Zöllner et al. 2024; Marleau et al. 2024).

Besides their large sizes and low surface brightnesses, many of these galaxies host surprisingly large globular cluster (GC) populations (van Dokkum et al. 2017, 2018b; Danieli et al.

2022), although some nearby dwarf irregular galaxies of similar luminosities have comparably rich GC populations (Georgiev et al. 2010; Forbes et al. 2018a, 2020).

Perhaps even more surprisingly, van Dokkum et al. (2018a) reported velocity dispersion measurement in the UDG NGC1052-DF2 (hereafter DF2) and argued that the total mass estimate using this measurement is consistent with the stellar mass of the galaxy. This indicated that the galaxy is not dark matter dominated, as is usual for galaxies of similar luminosity. van Dokkum et al. (2019a) reported a similar conclusion for the neighboring UDG NGC1052-DF4 (hereafter DF4) within the same group. Both of these galaxies have a sizeable globular cluster population with unusual GC luminosity function (van Dokkum et al. 2018b), and their GC luminosity function contains a strong peak at bright luminosities of $M_V \approx -9$ in addition to the usual peak at $M_V \sim -7.5$ (Shen et al. 2021a).

The conclusion about low dark matter content of DF2 and DF4 has been disputed. Martin et al. (2018) and Laporte et al. (2019) pointed out that uncertainties in velocity dispersion from globular clusters and luminosity measurement are substantial and the mass-to-light ratio for DF2 UDG is not constrained sufficiently well to exclude the presence of dark matter. They also argue that this galaxy is consistent with the scaling relations of Local Group dwarfs. Laporte et al. (2019) also pointed out that the dynamical mass estimated from velocity dispersion is a lower limit if there is significant rotation in the stellar or GC system. Indications of rotation have been reported for DF2 Lewis et al. (2020).

However, subsequent measurements presented by Danieli et al. (2019), Emsellem et al. (2019) and Shen et al. (2023) measured the velocity dispersion of DF2 and DF4 for stars and planetary nebulae, which indicated dispersion values consis-

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tent with the low values measured from globular clusters and thereby their low dark matter content. They also did not find indications of significant rotation. Nevertheless, the uncertainty of velocity dispersion measurements remains large. For example, [Emsellem et al. \(2019\)](#) report that the 95% confidence range after taking into account statistical and systematic uncertainties for DF2 galaxy is broad ($[0, 21]$ km/s).

Significant uncertainties and biases in the mass inference exist also (e.g., [Read et al. 2016b](#); [Oman et al. 2019](#); [Roper et al. 2023](#); [Downing & Oman 2023](#); [Sands et al. 2024](#)) for the dark matter mass inference of gas-rich field UDGs based on the mass modeling of their HI velocity fields [Guo et al. \(2020\)](#); [Mancera Piña et al. \(2020\)](#). As I will discuss below, the mass modeling uncertainties from the rotation curve measurement for the Small Magellanic Cloud (SMC) allow the SMC to be considered a dark matter-deficient galaxy.

As discussed in Section 4 below, several scenarios have been proposed and discussed both for the origin of UDGs and the origin of their dark matter deficient subset. Formation scenarios for the latter often invoke special processes, such as galaxy collisions or formation out of gas stripped from galaxies as a result of tidal interactions or ram pressure stripping. The baseline assumption of such scenarios is that the dark matter content of the UDGs is inconsistent with the expectation of regular galaxy formation.

In this study, I re-visit the question of whether constraints on the dark matter content of the UDGs, including the dark matter deficient ones, are consistent with expectations of the galaxy formation in Λ CDM. Specifically, I use the galaxy formation model of [Kravtsov & Manwadkar \(2022\)](#), which was demonstrated to reproduce properties of $L \lesssim L_\star$ galaxies down to ultra-faint dwarf luminosities at $z = 0$ ([Kravtsov & Manwadkar 2022](#); [Manwadkar & Kravtsov 2022](#); [Kravtsov & Wu 2023](#)), as well as properties of galaxies at $z > 5$ ([Kravtsov & Belokurov 2024](#); [Wu & Kravtsov 2024](#)).

The paper is organized as follows. The samples of observed galaxies used for comparisons with model results, the halo samples used to compute the evolution of galaxies in our model, and the galaxy formation itself are presented in Section 2. The main results of this study are presented in Section 3, where it is shown that the dark matter content of UDGs is consistent with the expectations of the model and that UDGs follow the same relation between dark matter mass within the half-mass radius and stellar mass as the dwarf satellites of the Milky Way. These results are discussed in the context of previous studies of UDGs and their proposed formation scenarios in Section 4. The main conclusions of this study are summarized in Section 5.

2. METHODS

2.1. Observational galaxy sample

In this study galaxy formation model results will be compared to the observed dwarf satellites of the Milky Way using V -band luminosities, half-light radii, and velocity dispersion measurements in Table B1 of [Kravtsov & Wu 2023](#), where measurements provenance is also specified. To convert V -band luminosities of satellites to stellar mass I use the M_\star/L_V predicted by the galaxy formation model, as quantified in [Kravtsov & Manwadkar \(2022\)](#): $M_\star = (2.85 + 0.1M_V)L_V$, where M_V is galaxy V -band absolute magnitude and its corresponding luminosity in solar luminosities is $L_V = 10^{0.4(M_{V\odot} - M_V)}$ and $M_{V\odot} = 4.81$ is the V -band absolute magnitude of the Sun ([Willmer 2018](#)).

In addition, I use a compilation of ultra-diffuse galaxies of [Gannon et al. \(2024\)](#).¹ Only galaxies with stellar velocity dispersion measurements are considered in this study. This is because the velocity dispersion and half-number radii for globular clusters can provide inaccurate or biased mass estimates (e.g., [Martin et al. 2018](#)). The ultra-diffuse galaxies are distinguished by the number of globular clusters into GC-rich and GC-poor systems. The GC-rich systems have more than 20 globular clusters and are listed in Table 2 of [Forbes & Gannon \(2024\)](#), I exclude VCC1448 because it does not meet UDG definition), the GC-poor UDGs are all other UDGs in the [Gannon et al. \(2024\)](#) compilation with measured stellar velocity dispersions. The stellar masses and effective radii used here are from this compilation. In addition, I use measurements for two DM-deficient galaxies NGC1052-DF2 ([van Dokkum et al. 2018a](#); [Danieli et al. 2019](#)) and NGC1052-DF4 ([van Dokkum et al. 2019a](#); [Shen et al. 2023](#)). Hereafter, the latter two galaxies will be referred to as DF2 and DF4 for brevity. For these galaxies I adopt stellar masses of $2 \times 10^8 M_\odot$ and $1.5 \times 10^8 M_\odot$, the projected circularized half-light radii of 2 and 1.6 kpc, and $M_{\text{tot}}(< r_{1/2})$ of $1.3 \pm 0.8 \times 10^8 M_\odot$ and $8^{+6}_{-4} \times 10^7 M_\odot$, respectively ([Danieli et al. 2019](#); [Shen et al. 2023](#)).

Finally, in the comparisons of the stellar mass-size relation of model and observed dwarf galaxies I use the sample of dwarf satellites in the ELVES survey (see Table 9 in [Carlsten et al. 2022](#)).

2.2. Halo tracks and catalogs

I model a population of galaxies in halos of virial mass $M_{200c} \lesssim 10^{12} M_\odot$ using halo evolutionary tracks from several suites of high-resolution N -body simulations of zoom-in regions around MW-sized halos. Specifically, I use the Caterpillar ([Griffen et al. 2016](#)) suite of N -body simulations² of 32 MW-sized haloes at the highest resolution level LX14 to maximize the dynamic range of halo masses probed by our modelling. The halo masses resolved in these simulations is sufficient to model the full range of luminosities of observed Milky Way satellites, as faintest ultrafaint dwarfs are hosted in haloes of $M_{\text{peak}} \gtrsim 10^7 M_\odot$ in our model ([Manwadkar & Kravtsov 2022](#)). The Caterpillar suite was simulated assuming flat Λ CDM cosmology with the Hubble constant of $h = H_0/100 = 0.6711$, the mean dimensionless matter density of $\Omega_{m0} = 0.32$, the rms fluctuations of matter density at the $8h^{-1}$ Mpc scale of $\sigma_8 = 0.8344$ and primordial power spectrum slope $n_s = 0.9624$.

In addition, I use halo tracks and catalogs from the ELVIS ([Garrison-Kimmel et al. 2014](#)) and Phat ELVIS (with central

¹ The compilation is available at https://github.com/gannonjs/Published_Data/tree/main/UDG_Spectroscopic_Data. Individual measurements are from [McConnachie \(2012\)](#); [van Dokkum et al. \(2015a\)](#); [Beasley et al. \(2016\)](#); [Martin et al. \(2016\)](#); [Yagi et al. \(2016b\)](#); [Martínez-Delgado et al. \(2016\)](#); [van Dokkum et al. \(2016, 2017\)](#); [Karachentsev et al. \(2017\)](#); [van Dokkum et al. \(2018a\)](#); [Toloba et al. \(2018\)](#); [Gu et al. \(2018\)](#); [Lim et al. \(2018\)](#); [Ruiz-Lara et al. \(2018\)](#); [Alabi et al. \(2018\)](#); [Ferré-Mateu et al. \(2018\)](#); [Forbes et al. \(2018b\)](#); [Martín-Navarro et al. \(2019\)](#); [Chilingarian et al. \(2019\)](#); [Fensch et al. \(2019\)](#); [Danieli et al. \(2019\)](#); [van Dokkum et al. \(2019b\)](#); [Torrealba et al. \(2019\)](#); [Iodice et al. \(2020\)](#); [Collins et al. \(2020\)](#); [Müller et al. \(2020\)](#); [Gannon et al. \(2020\)](#); [Lim et al. \(2020\)](#); [Müller et al. \(2021\)](#); [Forbes et al. \(2021\)](#); [Shen et al. \(2021b\)](#); [Gannon et al. \(2021, 2022b\)](#); [Mihos et al. \(2022\)](#); [Danieli et al. \(2022\)](#); [Villaume et al. \(2022\)](#); [Webb et al. \(2022\)](#); [Saifollahi et al. \(2022\)](#); [Janssens et al. \(2022\)](#); [Gannon et al. \(2023\)](#); [Ferré-Mateu et al. \(2023\)](#); [Toloba et al. \(2023\)](#); [Iodice et al. \(2023\)](#); [Shen et al. \(2023\)](#).

² <https://www.caterpillarproject.org>

disk and without Kelley et al. 2019) suites of zoom-in simulations of MW-sized halos. The mass and force resolution of these simulations are somewhat lower than those of the Caterpillar simulations (except for the two HiRes halos in the ELVIS suite). I use halos of mass $M_{200c} > 10^8 M_\odot$ from these simulations to complement halos from the Caterpillar suite to improve statistics of dwarf galaxies with stellar masses $M_\star > 10^5 M_\odot$. The ELVIS simulations were run assuming flat Λ CDM model with $h = H_0/100 = 0.71$, $\Omega_{m0} = 0.266$, $\sigma_8 = 0.801$, and $n_s = 0.963$, while Phat ELVIS simulations assumed $h = 0.6751$, $\Omega_{m0} = 0.3121$, $\sigma_8 = 0.81$, and $n_s = 0.968$.

I use evolution tracks of halos and subhalos that exist at $z = 0$ in and around the Milky Way-sized halo within zoom-in region in these simulation suites to construct mass evolution tracks for the galaxy formation model described below. The tracks were extracted from simulations using the Consistent Trees code (Behroozi et al. 2013) and consist of several halo properties, such as its virial mass, scale radius, maximum circular velocity, etc., measured at a series of redshifts from the first epoch at which progenitors are identified to $z = 0$. I use properties of dark matter halos at $z = 0$ to model the mass distribution in these halos.

Although some of the UDGs may result from tidal stripping, for many of them deep imaging does not show any signs of tidal interaction. Thus, in this analysis, observed UDGs are compared to the population of galaxies in halos that are not strongly affected by tidal stripping. To this end, I only consider halos with the ratio of the current mass with the peak virial mass along the track of $M_{200c}/M_{200c,peak} > 0.5$.

2.3. Galaxy formation model

The mass evolution tracks of halos in the zoom-in regions of the simulations described above were used as input for the GRUMPY galaxy formation model (Kravtsov & Manwadkar 2022). This is regulator-type galaxy formation model (e.g., Lilly et al. 2013) based on the models of Krumholz & Dekel (2012) and Feldmann (2013), but with a number of modifications to extend the model into the dwarf galaxy regime. The model evolves key properties of gas and stars (masses, metallicities, sizes, star formation rates, etc.) of the galaxies they host by solving a system of coupled differential equations governing the evolution of these properties. The model accounts for the UV heating after reionization and associated gas accretion suppression onto small mass haloes, galactic outflows, a model for gaseous disk and its size, molecular hydrogen mass, star formation, etc. The evolution of the half-mass radius of the stellar distribution is also modeled. The galaxy model parameters used in this study are identical to those used in Manwadkar & Kravtsov (2022).

The GRUMPY model is described and tested against a wide range of observations of local dwarf galaxies in Kravtsov & Manwadkar (2022). The model was shown to reproduce luminosity function and radial distribution of the Milky Way satellites and size-luminosity relation of observed dwarf galaxies (Manwadkar & Kravtsov 2022).

Figure 1 shows the relation between galaxy’s stellar mass and its stellar half-mass radius for model galaxies and observed galaxies. The latter include Milky Way satellites, the satellite galaxies in the ELVES sample (Carlsten et al. 2022), and the ultra-diffuse galaxies with existing measurements of stellar velocity dispersion (see Section 2.1 for details and references). Figure 1 shows that the median galaxy size–stellar mass relation and the scatter around it are in good agreement with the

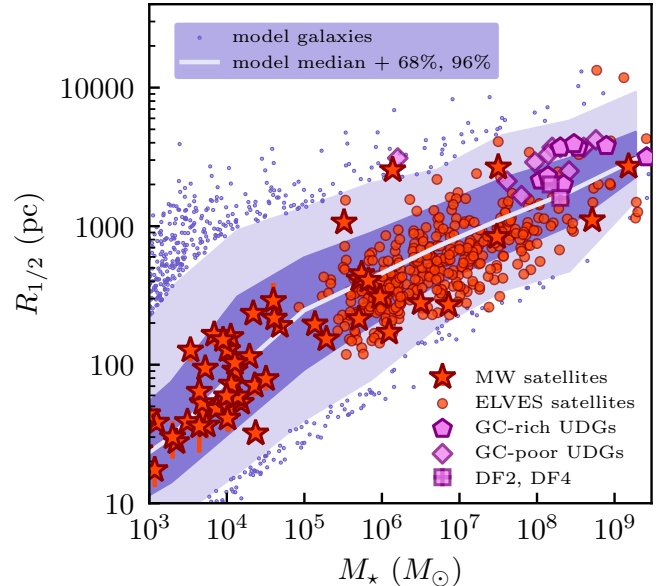


FIG. 1.— Half-mass radius of observed and model galaxies as a function of their stellar mass. Observed galaxies (see Sec. 2.1 for details and references) include Milky Way satellites (*red stars*), dwarf satellites in the ELVES host galaxy sample (*small circles*), GC-poor (*diamonds*) and GC-rich (*pentagons*) ultra-diffuse galaxies, and two DM-deficient galaxies NGC1052-DF2 and NGC1052-DF4. The median half-mass radius in bins of stellar mass for model galaxies is shown by the white line, while shaded bands show the 16th and 96th percentiles of the distribution around the median; blue dots show individual galaxies outside these ranges. The figure shows that the model reproduces $M_\star - R_{1/2}$ relation of observed galaxies and its scatter.

relation of observed dwarf galaxies (cf. also Manwadkar & Kravtsov 2022).

Predictably, given their definition, UDGs have sizes close to or above the median of the relation. I can also note that the DM-deficient galaxies DF2 and DF4 have some of the smallest sizes among the UDG sample shown in this figure. The UDG with $M_\star \approx 1.6 \times 10^6 M_\odot$ with a half-light radius just above the 96% band (farthest from the median among UDGs) is M31 dwarf spheroidal satellite And XIX (Martin et al. 2016). The Antlia II satellite of the Milky Way (Torrealba et al. 2019) is located near this galaxy. There are also several outliers in the ELVES satellite sample of Carlsten et al. (2022). This illustrates that ultra-diffuse galaxies exist and are fairly common among dwarf satellites of L_\star galaxies.

Note that the stellar mass-size relation is quite shallow at $M_\star > 10^5 M_\odot$, where size changes only by an order of magnitude over four orders of magnitude of stellar mass. The ultra-diffuse galaxies, which have $R_{1/2} > 1.5$ kpc by definition, thus span a wide range of stellar masses and, likely, halo masses. This is consistent with the finding of Danieli & van Dokkum (2019) that UDGs with a given size range span ≈ 6 absolute magnitudes or ≈ 2.5 orders of magnitude in luminosity.

In what follows, I use stellar masses and half-mass sizes produced by the model along with the information about halo mass and concentration of their parent halos from the halo catalogs, to estimate total and dark matter mass within the stellar half-mass radius, as I describe in the next section.

2.4. Mass profiles of model galaxies

To estimate total masses $M_{\text{tot}}(< r_{1/2})$ for model galaxies, I use individual half-mass radii, $r_{1/2}$, and consider two assump-

tions for the density profile of dark matter. Specifically, in the first model I adopt the Navarro-Frenk-White density profile (NFW, Navarro et al. 1997) for dark matter profile and use M_{200c} , R_{200c} , and the NFW scale radius, r_s , available in the halo catalogs:³

$$M_{\text{tot,NFW}}(< r_{1/2}) = M_{\text{dm,NFW}}(< r_{1/2}) + \frac{1}{2} M_{\star} \quad (1)$$

where M_{\star} is the stellar mass of the galaxy hosted by a given halo computed in the model and

$$M_{\text{dm,NFW}}(< r) = M_{200c} \frac{f(r/r_s)}{f(R_{200c}/r_s)}; \quad (2)$$

where

$$f(x) = \ln(1+x) - \frac{x}{x+1}. \quad (3)$$

In the second model, I assume the dark matter density profile of Lazar et al. (2020), which approximates the effects of stellar feedback on the density profile in the FIRE-2 simulations:

$$M_{\text{tot,L}}(< r_{1/2}) = M_{\text{dm,Lazar}}(< r_{1/2}) + \frac{1}{2} M_{\star}. \quad (4)$$

Specifically, I use the parametrization of the cored-Einasto density profile in the equations 8-10, 12 of Lazar et al. (2020) and equations for the cumulative mass profile in their Appendix B1 and parameters in the second row of their Table 1 for the dependence of profile as a function of stellar mass M_{\star} . I chose to use the dependence on the stellar mass, to minimize the effects of different M_{\star}/M_h in their simulations and our model. Note that the profiles were calibrated only for galaxies of $M_{\star} \gtrsim 10^5 M_{\odot}$. However, in this model effects of feedback for $M_{\star} < 10^6 M_{\odot}$ are expected to be negligible and thus extrapolating their results to smaller masses is equivalent to simply assuming Einasto profile with negligible core for these low-mass systems.

Note that I neglect the gas mass in the equations 1 and 4. Given that most of the observed UDGs I compare with are located in galaxy groups and clusters I assume that most of their ISM gas is stripped.

2.5. Stellar mass–halo mass relation of model galaxies

Figure 2 shows the relation between the stellar mass of model galaxies and the peak virial mass of their host halos for the stellar mass range comparable to that of the observed galaxy sample that the model is compared to in this study. The peak mass is estimated as the maximum virial mass along the halo evolution track. The physical origin of this relation and its shape are discussed in Kravtsov & Manwadkar (2022) and Manwadkar & Kravtsov (2022) and I refer the interested reader to these papers for details. I note that this relation agrees with existing direct observational constraints on stellar and halo masses of galaxies in this mass range and with the luminosity function of Milky Way satellites, as well as with the relations in the current state-of-the-art cosmological simulations of dwarf galaxy formation (Manwadkar & Kravtsov 2022).

Here I reproduce the relation for the model galaxy samples used in this study to illustrate the halo masses in which model galaxies of a given stellar mass reside. In particular, most of the UDGs I will examine have stellar masses of $M_{\star} \approx$

³ Note that for subhalos the scale radius and M_{200c} are estimated using only bound dark matter particles by the Rockstar halo finder (Behroozi et al. 2013).

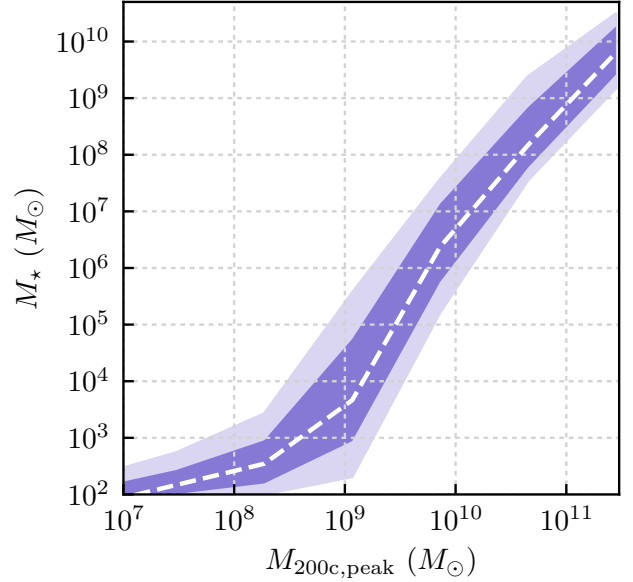


FIG. 2.— Stellar mass of model galaxies as a function of peak virial mass, $M_{200c,\text{peak}}$ of their host halos. The dashed line shows the median stellar mass in bins of halo mass, while the dark and light shaded bands show the 58th and 96th percentiles of M_{\star} distribution.

$5 \times 10^7 - 10^9 M_{\odot}$, which corresponds to halo masses of $M_{200c} \approx 8 \times 10^9 - 2 \times 10^{11} M_{\odot}$. This is broadly consistent with the constraints on UDG halo masses of $M_{200c} \leq 10^{11.8} M_{\odot}$ using stacked weak lensing signal (Sifón et al. 2018), estimates of their halo masses using halo mass-globular cluster relation (Prole et al. 2019). This is also consistent with the halo mass range $M_{200c} < 10^{11.2} M_{\odot}$ of model UDGs in the Illustris TNG (Benavides et al. 2023; Doppel et al. 2024).

I will also compare these galaxies to dwarf satellite galaxies of the Milky Way, which span a much wider range of masses down to $M_{\star} \approx 10^3 M_{\odot}$ or, according to Figure 2, halos of masses $M_{200c} \approx 10^8 - 10^9 M_{\odot}$. Comparisons of the model $M_{\text{tot}}(< r_{1/2})$ and $M_{\text{dm}}(< r_{1/2})$ with the values estimated for observed galaxies of the same stellar mass will serve as a test of whether model forms galaxies in halos of correct mass and concentration.

3. RESULTS

Figure 3 shows the total mass within half-light radius for model galaxies and observational estimates for the MW dwarf satellites, DM-deficient UDGs DF2 and DF4, and the sample of UDGs with measured stellar velocity dispersions compiled by Gannon et al. (2024). The figure shows a remarkable agreement between the Λ CDM-based galaxy formation model results (lines and shaded bands) and the relation traced by the observed galaxies shown as points. The figure also shows that UDGs follow the relation of the Milky Way dwarf satellites.

The UDGs are also broadly consistent with model expectations, but I note that quite a few are located outside the 1σ model band. On the low side, DF2, DF4 and PUDG-R15 galaxies are at $\approx -1\sigma$ from the model median. On the high side, DGSAT-1 UDG is outside 2σ band and DF44 is close to it. The Small Magellanic Cloud (SMC) galaxy has comparable baryon and dark matter masses within half-light radius (see below), while the Nube galaxy with a similar stellar mass and half-light radius of 6.9 kpc is dark matter dominated (Montes et al. 2024).

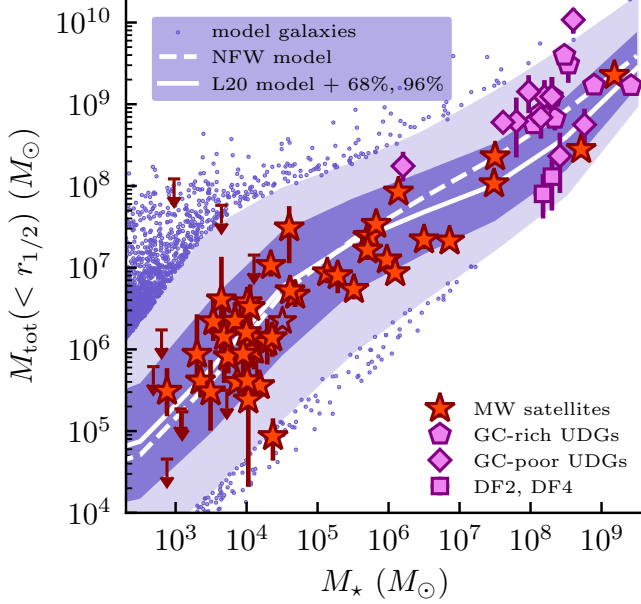


FIG. 3.— Total mass within half-light radius vs stellar mass of model galaxies and dwarf satellites of the Milky Way (stars and upper limits) and UDGs (the sample shown in Figure 1). The dashed and solid lines show the median relations for model galaxies for $M_{\text{tot}}(< r_{1/2})$ computed assuming NFW and Lazar et al. (2020) dark matter density profiles, respectively. The dark- and light-shaded bands show the 16th and 96th percentiles of the distribution around the median of the latter model; blue dots show individual galaxies outside these ranges. The model matches the $M_{\star} - M_{\text{tot}}(< r_{1/2})$ relation well at all stellar masses. The UDGs tend to lie above the median model relation. The DM-deficient galaxies DF2 and DF4 lie below the median relation but within the 96% of the model galaxy distribution for their stellar mass.

One can note that $M_{\text{tot}}(< r_{1/2})$ becomes close to M_{\star} for galaxies with $M_{\star} > 10^8 M_{\odot}$. This is due to a general trend of decreasing $M_{\text{tot}}(< r_{1/2})$ -to-luminosity ratio with increasing luminosity for dwarf galaxies (e.g., Wolf et al. 2010). Nevertheless, each model galaxy has dark matter within $r_{1/2}$ and thus $M_{\text{tot}}(< r_{1/2}) > M_{\star}/2$. Figure 4 shows $M_{\text{dm}}(< r_{1/2}) = M_{\text{tot}}(< r_{1/2}) - M_{\star}/2$ as a function of M_{\star} . In most observed galaxies $M_{\text{dm}}(< r_{1/2}) - \sigma_{M_{\text{dm}}} > 0$ and their $M_{\text{dm}}(< r_{1/2})$ are shown as points with errorbars. Galaxies with $M_{\text{dm}}(< r_{1/2}) - \sigma_{M_{\text{dm}}} < 0$ are shown as upper limits with the top of the limit corresponding to $M_{\text{dm}}(< r_{1/2}) + \sigma_{M_{\text{dm}}}$.

The figure shows that there are only 4 such galaxies: three UDGs – DF2, DF4 (“dark matter-deficient” UDGs in the NGC 1052 group, Danieli et al. 2019; Shen et al. 2023) and PUDG-R15 (a UDG in the Perseus cluster Gannon et al. 2022a) – and SMC.⁴ In all four of these galaxies $M_{\text{dm}}(< r_{1/2})$ is consistent with zero. However, Figure 4 shows that the upper limits on $M_{\text{dm}}(< r_{1/2})$ in these galaxies are also consistent with expectation of the Λ CDM-based galaxy formation model. In other words, the amount of the expected DM within half-light radii of these galaxies is consistent with observational upper limits.

Galaxies of a given stellar mass have a broad range of sizes which translates to a broad range of $M_{\text{tot}}(< r_{1/2})$ at

⁴ The SMC has the half-light radius of 1106 ± 77 pc (Muñoz et al. 2018) and stellar mass within this radius of $1.5 - 2.5 \times 10^8 M_{\odot}$ (given the estimates of total stellar mass of $3 - 5 \times 10^8 M_{\odot}$ Harris & Zaritsky 2004; Skibba et al. 2012; Di Teodoro et al. 2019), while its rotation velocity of 33 ± 2 km/s (Di Teodoro et al. 2019) at the half-light radius implies $M_{\text{tot}}(< r_{1/2}) = 2.8 \pm 0.4 \times 10^8 M_{\odot}$.

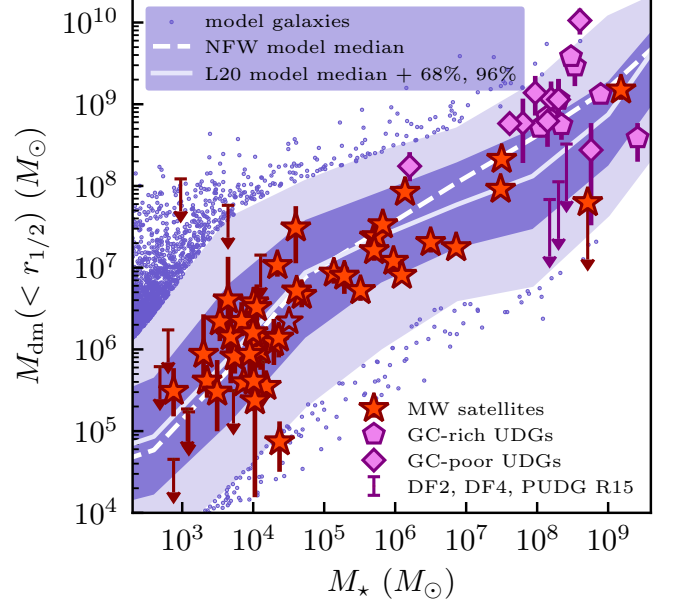


FIG. 4.— Dark matter mass within half-light radius, $M_{\text{dm}}(< r_{1/2}) = M_{\text{tot}}(< r_{1/2}) - M_{\star}/2$ vs stellar mass of model galaxies and dwarf satellites of the Milky Way (stars and light red upper limits) and UDGs (squares, diamonds, pentagons, and dark red upper limits). $M_{\text{dm}}(< r_{1/2})$ is computed as $M_{\text{tot}}(< r_{1/2}) - M_{\star}/2$. The galaxies for which $M_{\text{dm}}(< r_{1/2}) - \sigma_{M_{\text{dm}}} < 0$ are shown as upper limits with the top of the limit corresponding to $M_{\text{dm}}(< r_{1/2}) + \sigma_{M_{\text{dm}}}$. The dashed and solid lines show the median relations for model galaxies for $M_{\text{tot}}(< r_{1/2})$ computed assuming NFW and Lazar et al. (2020) dark matter density profiles, respectively. The dark- and light-shaded bands show the 16th and 96th percentiles of the distribution around the median of the latter model; blue dots show individual galaxies outside these ranges. The model matches the $M_{\star} - M_{\text{dm}}(< r_{1/2})$ relation followed by observed galaxies well at all stellar masses. The UDGs tend to lie above the median model relation. The DM-deficient galaxies DF2 and DF4 are shown as upper limits as their $M_{\text{dm}}(< r_{1/2})$ is consistent with zero within its uncertainty. This is also true for PUDG-R15 UDG from the sample of Gannon et al. (2022a), shown as the third dark red upper limit, and for SMC shown as the star with light red upper limit. However, the figure shows that these upper limits are fully consistent with the expected distribution of $M_{\text{dm}}(< r_{1/2})$ at the stellar mass of these galaxies. They are thus not DM-deficient in the context of this model.

a given stellar mass (as also noted for observed UDGs by van Dokkum et al. 2019b). Given that $M_{\text{tot}}(< r_{1/2})$ depends on $r_{1/2}$, one may wonder if the model galaxies that have $M_{\text{dm}}(< r_{1/2})$ consistent with the constraints of DF2, DF4, and PUDG-R15 have predominantly sizes smaller than the sizes of these galaxies. Figure 5 shows that this is not the case. This figure compares $M_{\text{dm}}(< r_{1/2})$ as a function of $R_{1/2}$ for observed and model galaxies in the stellar mass range of $7 \times 10^7 < M_{\star}/M_{\odot} < 3 \times 10^8 M_{\odot}$. The dark matter mass for model galaxies is shown for both the NFW and Lazar et al. (2020) dark matter density profiles. Figure 5 shows that the upper limits for DF2, DF4, and PUDG-R15 galaxies are consistent with $M_{\text{dm}}(< r_{1/2})$ in model galaxies of comparable size, even though the upper limits on dark matter mass for these observed galaxies were estimated without taking into account uncertainty of their stellar mass estimates. Other UDGs are also consistent with model expectations and lie on the $M_{\text{dm}} - R_{1/2}$ relation traced by model galaxies.

The overall dependence of $M_{\text{dm}}(< r_{1/2})$ on $r_{1/2}$ reflects the dark matter mass profile of halos hosting galaxies of this stellar mass. The effect of feedback on the mass profile at this stellar mass ($M_{\star} \approx 10^8 M_{\odot}$) quantified by the Lazar et al.

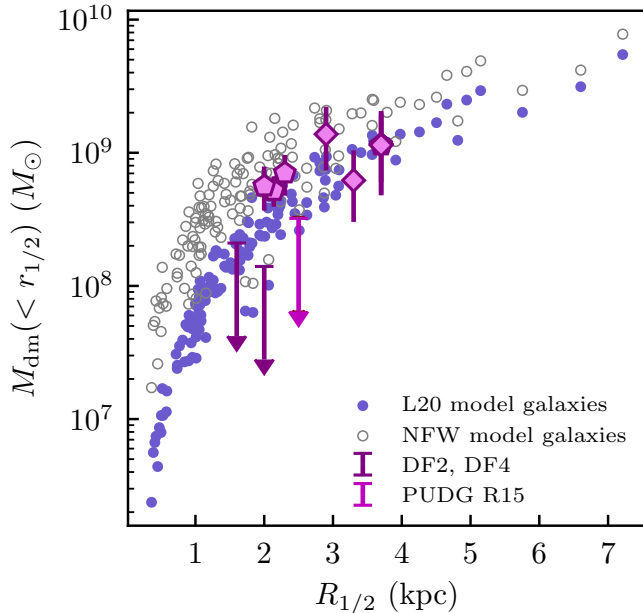


FIG. 5.— Dark matter mass within half-light radius, $M_{\text{dm}}(< r_{1/2})$ vs half-light radius for model galaxies in the stellar mass range of $7 \times 10^7 < M_{\star}/M_{\odot} < 3 \times 10^8 M_{\odot}$ assuming NFW (open gray circles) and Lazar et al. (2020, blue circles) dark matter density profiles. The upper limits show constraints on $M_{\text{dm}}(< r_{1/2})$ for observed galaxies DF2, DF4 and PUDG-R15, while other UDGs in this mass range are shown as points with the same type as in Figures 3 and 4.

(2020) profile is larger than the scatter of M_{dm} at a given $R_{1/2}$. This effect brings model masses in better agreement with observational estimates of $M_{\text{tot}}(< r_{1/2})$ for UDGs, although even unmodified NFW mass profiles are still fairly consistent with these estimates.

The scatter of M_{dm} for a given value of $R_{1/2}$ for the Lazar et al. (2020) mass profile is due to 1) scatter of halo concentrations and 2) scatter of M_{\star}/M_{200c} due to the dependence of the Lazar et al. (2020) profile on this ratio. For the galaxies of $M_{\star} \sim 10^8 M_{\odot}$ with sizes of $1.5 < r_{1/2} < 3$ kpc shown in Figure 5, there is a positive correlation between $M_{\text{dm}}(< r_{1/2})$ and concentration c_{200c} with the correlation coefficient of 0.51 ± 0.09 , where uncertainty is estimated using bootstrap resampling. There is also a weaker anti-correlation between $M_{\text{dm}}(< r_{1/2})$ and $\log_{10} M_{\star}/M_{200c}$ with the correlation coefficient of -0.31 ± 0.11 .

Thus, in the galaxy formation model used here, the most dark matter-deficient galaxies of a given size correspond to halos with the smallest concentrations and the largest ratios of M_{\star}/M_{200c} . Conversely, the most dark matter-dominated galaxies are hosted by the highest concentration halos with the smallest M_{\star}/M_{200c} ratios. Note that halo concentration correlates with halo formation time with low (high) concentration halos forming late (early). The M_{\star}/M_{200c} ratio depends on the interplay between gas accretion, star formation, and feedback-driven outflows for different halo mass assembly histories.

4. DISCUSSION

The main conclusion that can be drawn from the results presented above is that the dark matter content of UDGs is broadly consistent with the expectation of our Λ CDM-based galaxy formation model. The large variation in the amount of dark matter contribution to the total mass within $r_{1/2}$ for observed

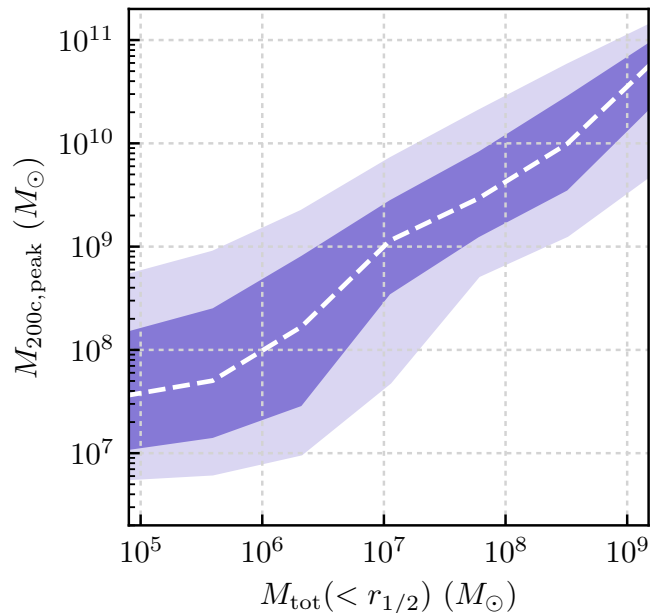


FIG. 6.— Relation between total mass within half-light radius, $M_{\text{tot}}(< r_{1/2})$, and the peak virial mass of the halo, $M_{200c,\text{peak}}$. The dashed line shows the median of the distribution of $M_{200c,\text{peak}}$ in bins of $M_{\text{tot}}(< r_{1/2})$, while dark- and light-shaded bands show 68th and 96th percentiles of its distribution. The figure shows that a well-defined relation between $M_{\text{tot}}(< r_{1/2})$ and $M_{200c,\text{peak}}$ exists for model galaxies, but the scatter around the median relation is significant.

dwarf galaxies and UDGs of the same stellar mass – from very dark matter-dominated galaxies DGSAT-1 and DF44 to “dark matter-deficient” galaxies DF2, DF4, and PUDG-R15 – is due mainly to 1) the large range of sizes of galaxies of a given M_{\star} (see Figure 1), 2) the concentration of their halo and 3) feedback effects expected for the galaxies of such stellar mass (e.g., Governato et al. 2012; Pontzen & Governato 2012; Di Cintio et al. 2014; Chan et al. 2015; Tollet et al. 2016; Read et al. 2016a; Lazar et al. 2020; Brook et al. 2021). This variation is thus much larger than the expected scatter in the dark matter density profiles at a fixed radius for galaxies of a given stellar mass (as is the case in observed UDGs van Dokkum et al. 2019b).

For these reasons, the range of virial halo masses for a given $M_{\text{tot}}(< r_{1/2})$ is also quite broad. Figure 6 shows the peak virial mass, $M_{200c,\text{peak}}$ of host halos hosting model galaxies for galaxies with a given $M_{\text{tot}}(< r_{1/2})$. The median of the distribution is shown by the dashed line, while 68th and 96th percentiles are shown by the dark and light shaded areas. Although there is a clear relationship between the two masses, the 96th percentile range of the virial mass spans ≥ 1.5 dex at a given $M_{\text{tot}}(< r_{1/2})$. Specifically, galaxies with $M_{\text{tot}}(< r_{1/2}) = 5 \times 10^8 M_{\odot}$ typical for the UDGs shown in the figures can be hosted in halos of M_{200c} as low as $\approx 3 \times 10^9 M_{\odot}$ and as high as $\approx 8 \times 10^{10} M_{\odot}$. Still, the halo mass range of observed UDGs indicated by our model implies that these galaxies are hosted by halos of sub-LMC mass, consistent with conclusions of Rong et al. (2017) based on the analyses of the Millenium simulation coupled with a semi-analytic galaxy formation model. This is also consistent with the lack of X-rays from hot gas and X-ray binaries from UDGs that were believed to be most massive (Kovács et al. 2019; Bogdán 2020).

The origin of UDGs, and especially the origin of the dark

matter-deficient UDGs, is still actively debated. One class of models for abundant UDGs in cluster environment considers ram pressure stripping of gas, accompanied by subsequent suppression of gas accretion by the intracluster medium and tidal heating (Yozin & Bekki 2015; Safarzadeh & Scannapieco 2017; Carleton et al. 2019; Tremmel et al. 2020; Sales et al. 2020; Moreno et al. 2022; Fielder et al. 2024). Another class invokes internal processes, such as stellar feedback (Di Cintio et al. 2017; Chan et al. 2018; Martin et al. 2019; Jiang et al. 2019; Jackson et al. 2021; Trujillo-Gomez et al. 2021, 2022) or inefficient cooling (Wang et al. 2023), leading to large sizes and low surface brightness and star formation. Wright et al. (2021) find that in their simulations field UDGs form via early mergers, which induce high spin onto the gas and temporarily boost and redistribute gas and star formation to the outskirts of galaxies, resulting in lower central SFRs and lower surface brightness later on. Finally, remnants of high-speed collisions of galaxies were proposed as a possible origin for dark-matter deficient UDGs (Baushev 2018) and this scenario is supported by simulation results (Lee et al. 2021, 2024). Van Nest et al. (2022) point out that the interpretation of the UDG formation may depend on the details of the UDG definition and may be different for diffuse galaxies in clusters and in the field. Furthermore, a combination of various processes may be at play during UDG formation (e.g., Jackson et al. 2021).

van Dokkum et al. (2022a) argued that DF2 and DF4 UDGs may have formed from gas stripped during a collision of a dwarf galaxy with a massive galaxy NGC 1052. The common formation process of these galaxies is indicated by the similarity of colors of their globular clusters (van Dokkum et al. 2022b). The tidal dwarf formation scenario for dark matter-deficient UDGs was also explored by Ivleva et al. (2024) but in their scenario UDGs form from the gas stripped from galaxies by the ram pressure from the intracluster medium. Keim et al. (2022, cf. also Golini et al. 2024) present evidence that DF2 and DF4 UDGs may undergo a significant tidal stripping, which may well have reduced the dark matter mass within their half-light radius. However, other studies found no signs of tidal interaction in DF4 even in very deep images (Montes et al. 2021; Golini et al. 2024).

Results presented in this paper do not exclude any of the specific formation scenarios discussed in the literature, but they imply that a special formation mechanism for DM-deficient UDGs is *not required* because their dark matter content constraints are consistent with the range of dark matter masses expected for halos hosting UDGs. This is akin to the overall interpretation of UDG population as a high-spin/large-size tail of the galaxy distribution (Amorisco & Loeb 2016, cf. also Dalcanton et al. 1997, Rong et al. 2017, 2024, Liao et al. 2019 and Benavides et al. 2023). Likewise, recent observational studies find that UDGs in clusters correspond structurally to the low-surface brightness tail of the regular dwarf galaxy population (Martin et al. 2018; Trujillo 2021; Zöller et al. 2024; Marleau et al. 2024).

5. CONCLUSIONS

I have compared the dark matter content within a half-mass radius in model galaxies formed within a realistic Λ CDM-based galaxy formation model and observed dwarf satellites of the Milky Way and ultra-diffuse galaxies. The model reproduces the main properties and scaling relations of dwarf galaxies, in particular their stellar mass-size relation (Manwadkar & Kravtsov 2022, see also Fig. 1 above). The main results and conclusions of this study can be summarized as

follows.

1. The model reproduces the total mass, $M_{\text{tot}}(< r_{1/2}) = M_{\star}/2 + M_{\text{dm}}(< r_{1/2})$, and dark matter mass, $M_{\text{dm}}(< r_{1/2})$, within half-mass radius of observed dwarf galaxies. It also reproduces well the correlation of these masses with stellar mass exhibited by observed galaxies (see Figures 3 and 4).
2. Ultra-diffuse galaxies lie on the same $M_{\text{dm}}(< r_{1/2}) - M_{\star}$ relation as the Milky Way dwarf satellites but their masses are generally larger than the median expected mass due to their large sizes (Figures 4 and 5).
3. The scatter in the $M_{\text{dm}}(< r_{1/2}) - M_{\star}$ relation is driven primarily by the broad range of sizes of galaxies of a given stellar mass (Figure 5). This large scatter results in a wide range of halo virial masses that may correspond to a given value of mass within $r_{1/2}$ (Figure 6).
4. The most dark matter-deficient galaxies of a given size correspond to halos with the smallest concentrations and the largest ratios of M_{\star}/M_{200c} . Conversely, the most dark matter-dominated galaxies are hosted by the highest concentration halos with the smallest M_{\star}/M_{200c} ratios.
5. The model indicates that UDGs with stellar mass of $M_{\star} \approx 10^8 M_{\odot}$ form in halos of the present-day virial mass $M_{200c} \approx 10^{10} - 10^{11} M_{\odot}$ (see Figure 2). The scatter between $M_{\text{dm}}(< r_{1/2})$ and M_{200c} , however, is expected to be large (Figure 6, which renders inference of the virial mass from $M_{\text{dm}}(< r_{1/2})$ uncertain and dependent on specific assumptions about the halo mass profile).

Results presented in this paper indicate that dark matter-deficient UDGs may represent a tail of the expected dark matter profiles, especially if the effect of feedback on these profiles is taken into account. It will be important to test this conclusion in the future using larger samples of ultra-diffuse galaxies.

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DATA AVAILABILITY

Halo catalogs from the Caterpillar simulations are available at <https://www.caterpillarproject.org/>. The GRUMPY model pipeline is available at <https://github.com/kibokov/GRUMPY>. The data used in the plots within this article are available on request to the author.

REFERENCES

- Alabi A., et al., 2018, *MNRAS*, **479**, 3308
- Amorisco N. C., Loeb A., 2016, *MNRAS*, **459**, L51
- Baushev A. N., 2018, *New A.*, **60**, 69
- Bautista J. M. G., Koda J., Yagi M., Komiyama Y., Yamanoi H., 2023, *ApJS*, **267**, 10
- Beasley M. A., Romanowsky A. J., Pota V., Navarro I. M., Martinez Delgado D., Neyer F., Deich A. L., 2016, *ApJ*, **819**, L20
- Behroozi P. S., Wechsler R. H., Wu H.-Y., Busha M. T., Klypin A. A., Primack J. R., 2013, *ApJ*, **763**, 18
- Benavides J. A., Sales L. V., Abadi M. G., Marinacci F., Vogelsberger M., Hernquist L., 2023, *MNRAS*, **522**, 1033
- Bogdán Á., 2020, *ApJ*, **901**, L30
- Bothun G. D., Beers T. C., Mould J. R., Huchra J. P., 1985, *AJ*, **90**, 2487
- Bothun G. D., Impey C. D., Malin D. F., Mould J. R., 1987, *AJ*, **94**, 23
- Bothun G., Impey C., McGaugh S., 1997, *PASP*, **109**, 745
- Brook C. B., Di Cintio A., Macciò A. V., Blank M., 2021, *ApJ*, **919**, L1
- Carleton T., Errani R., Cooper M., Kaplinghat M., Peñarrubia J., Guo Y., 2019, *MNRAS*, **485**, 382
- Carlsten S. G., Greene J. E., Beaton R. L., Danieli S., Greco J. P., 2022, *ApJ*, **933**, 47
- Chan T. K., Kereš D., Oñorbe J., Hopkins P. F., Muratov A. L., Faucher-Giguère C. A., Quataert E., 2015, *MNRAS*, **454**, 2981
- Chan T. K., Kereš D., Wetzel A., Hopkins P. F., Faucher-Giguère C. A., El-Badry K., Garrison-Kimmel S., Boylan-Kolchin M., 2018, *MNRAS*, **478**, 906
- Chilingarian I. V., Afanasiev A. V., Grishin K. A., Fabricant D., Moran S., 2019, *ApJ*, **884**, 79
- Collins M. L. M., Tollerud E. J., Rich R. M., Ibata R. A., Martin N. F., Chapman S. C., Gilbert K. M., Preston J., 2020, *MNRAS*, **491**, 3496
- Conselice C. J., Gallagher John S. I., Wyse R. F. G., 2003, *AJ*, **125**, 66
- Dalcanton J. J., Spergel D. N., Gunn J. E., Schmidt M., Schneider D. P., 1997, *AJ*, **114**, 635
- Danieli S., van Dokkum P., 2019, *ApJ*, **875**, 155
- Danieli S., van Dokkum P., Conroy C., Abraham R., Romanowsky A. J., 2019, *ApJ*, **874**, L12
- Danieli S., et al., 2022, *ApJ*, **927**, L28
- Di Cintio A., Brook C. B., Dutton A. A., Macciò A. V., Stinson G. S., Knebe A., 2014, *MNRAS*, **441**, 2986
- Di Cintio A., Brook C. B., Dutton A. A., Macciò A. V., Obreja A., Dekel A., 2017, *MNRAS*, **466**, L1
- Di Teodoro E. M., et al., 2019, *MNRAS*, **483**, 392
- Disney M. J., 1976, *Nature*, **263**, 573
- Doppel J. E., Sales L. V., Benavides J. A., Toloba E., Peng E. W., Nelson D., Navarro J. F., 2024, *MNRAS*, **529**, 1827
- Downing E. R., Oman K. A., 2023, *MNRAS*, **522**, 3318
- Driver S. P., Liske J., Cross N. J. G., De Propriis R., Allen P. D., 2005, *MNRAS*, **360**, 81
- Emsellem E., et al., 2019, *A&A*, **625**, A76
- Feldmann R., 2013, *MNRAS*, **433**, 1910
- Fensch J., et al., 2019, *A&A*, **625**, A77
- Ferré-Mateu A., et al., 2018, *MNRAS*, **479**, 4891
- Ferré-Mateu A., Gannon J. S., Forbes D. A., Buzzo M. L., Romanowsky A. J., Brodie J. P., 2023, *MNRAS*, **526**, 4735
- Fielder C., Jones M., Sand D., Bennet P., Crnojevic D., Karunakaran A., Mutlu-Pakdil B., Spekkens K., 2024, *arXiv e-prints*, p. arXiv:2401.01931
- Forbes D. A., Gannon J., 2024, *MNRAS*, **528**, 608
- Forbes D. A., Read J. I., Gieles M., Collins M. L. M., 2018a, *MNRAS*, **481**, 5592
- Forbes D. A., Read J. I., Gieles M., Collins M. L. M., 2018b, *MNRAS*, **481**, 5592
- Forbes D. A., Alabi A., Romanowsky A. J., Brodie J. P., Arimoto N., 2020, *MNRAS*, **492**, 4874
- Forbes D. A., Gannon J. S., Romanowsky A. J., Alabi A., Brodie J. P., Couch W. J., Ferré-Mateu A., 2021, *MNRAS*, **500**, 1279
- Forbes D. A., et al., 2023, *MNRAS*, **525**, L93
- Gannon J. S., Forbes D. A., Romanowsky A. J., Ferré-Mateu A., Couch W. J., Brodie J. P., 2020, *MNRAS*, **495**, 2582
- Gannon J. S., et al., 2021, *MNRAS*, **502**, 3144
- Gannon J. S., et al., 2022a, *MNRAS*, **510**, 946
- Gannon J. S., et al., 2022b, *MNRAS*, **510**, 946
- Gannon J. S., Forbes D. A., Brodie J. P., Romanowsky A. J., Couch W. J., Ferré-Mateu A., 2023, *MNRAS*, **518**, 3653
- Gannon J. S., Ferré-Mateu A., Forbes D. A., Brodie J. P., Buzzo M. L., Romanowsky A. J., 2024, *arXiv e-prints*, p. arXiv:2405.09104
- Garrison-Kimmel S., Boylan-Kolchin M., Bullock J. S., Lee K., 2014, *MNRAS*, **438**, 2578
- Geller M. J., Diaferio A., Kurtz M. J., Dell'Antonio I. P., Fabricant D. G., 2012, *AJ*, **143**, 102
- Georgiev I. Y., Puzia T. H., Goudfrooij P., Hilker M., 2010, *MNRAS*, **406**, 1967
- Golini G., Montes M., Carrasco E. R., Román J., Trujillo I., 2024, *A&A*, **684**, A99
- Governato F., et al., 2012, *MNRAS*, **422**, 1231
- Griffen B. F., Ji A. P., Dooley G. A., Gómez F. A., Vogelsberger M., O'Shea B. W., Frebel A., 2016, *ApJ*, **818**, 10
- Gu M., et al., 2018, *ApJ*, **859**, 37
- Guo Q., et al., 2020, *Nature Astronomy*, **4**, 246
- Harris J., Zaritsky D., 2004, *AJ*, **127**, 1531
- Harris C. R., et al., 2020, *Nature*, **585**, 357–362
- Hunter J. D., 2007, *Computing In Science & Engineering*, **9**, 90
- Impey C., Bothun G., 1997, *ARA&A*, **35**, 267
- Impey C., Bothun G., Malin D., 1988, *ApJ*, **330**, 634
- Iodice E., et al., 2020, *A&A*, **642**, A48
- Iodice E., et al., 2023, *arXiv e-prints*, p. arXiv:2308.11493
- Ivleva A., Remus R.-S., Valenzuela L. M., Dolag K., 2024, *arXiv e-prints*, p. arXiv:2402.09060
- Jackson R. A., et al., 2021, *MNRAS*, **502**, 4262
- Janssens S. R., et al., 2022, *MNRAS*, **517**, 858
- Jiang F., Dekel A., Freundlich J., Romanowsky A. J., Dutton A. A., Macciò A. V., Di Cintio A., 2019, *MNRAS*, **487**, 5272
- Karachentsev I. D., Makarova L. N., Sharina M. E., Karachentseva V. E., 2017, *Astrophysical Bulletin*, **72**, 376
- Keim M. A., et al., 2022, *ApJ*, **935**, 160
- Kelley T., Bullock J. S., Garrison-Kimmel S., Boylan-Kolchin M., Pawlowski M. S., Graus A. S., 2019, *MNRAS*, **487**, 4409
- Koda J., Yagi M., Yamanoi H., Komiyama Y., 2015, *ApJ*, **807**, L2
- Kovács O. E., Bogdán Á., Canning R. E. A., 2019, *ApJ*, **879**, L12
- Kravtsov A., Belokurov V., 2024, *arXiv e-prints*, p. arXiv:2405.04578
- Kravtsov A., Manwadkar V., 2022, *MNRAS*, **514**, 2667
- Kravtsov A., Wu Z., 2023, *MNRAS*, **525**, 325
- Krumholz M. R., Dekel A., 2012, *ApJ*, **753**, 16
- Laporte C. F. P., Agnello A., Navarro J. F., 2019, *MNRAS*, **484**, 245
- Lazar A., et al., 2020, *MNRAS*, **497**, 2393
- Lee J., Shin E.-j., Kim J.-h., 2021, *ApJ*, **917**, L15
- Lee J., Shin E.-j., Kim J.-h., Shapiro P. R., Chung E., 2024, *ApJ*, **966**, 72
- Leisman L., et al., 2017, *ApJ*, **842**, 133
- Lewis G. F., Brewer B. J., Wan Z., 2020, *MNRAS*, **491**, L1
- Liao S., et al., 2019, *MNRAS*, **490**, 5182
- Lilly S. J., Carollo C. M., Pipino A., Renzini A., Peng Y., 2013, *ApJ*, **772**, 119
- Lim S., Peng E. W., Côté P., Sales L. V., den Brok M., Blakeslee J. P., Guhathakurta P., 2018, *ApJ*, **862**, 82
- Lim S., et al., 2020, *ApJ*, **899**, 69
- Mancera Piña P. E., Peletier R. F., Aguerri J. A. L., Venhola A., Trager S., Choque Challapa N., 2018, *MNRAS*, **481**, 4381
- Mancera Piña P. E., Aguerri J. A. L., Peletier R. F., Venhola A., Trager S., Choque Challapa N., 2019, *MNRAS*, **485**, 1036
- Mancera Piña P. E., et al., 2020, *MNRAS*, **495**, 3636
- Manwadkar V., Kravtsov A. V., 2022, *MNRAS*, **516**, 3944
- Marleau F. R., et al., 2024, *arXiv e-prints*, p. arXiv:2405.13502
- Martín-Navarro I., et al., 2019, *MNRAS*, **484**, 3425
- Martin N. F., et al., 2016, *ApJ*, **833**, 167
- Martin N. F., Collins M. L. M., Longeard N., Tollerud E., 2018, *ApJ*, **859**, L5
- Martin G., et al., 2019, *MNRAS*, **485**, 796
- Martínez-Delgado D., et al., 2016, *AJ*, **151**, 96
- McConnachie A. W., 2012, *AJ*, **144**, 4
- McGaugh S. S., Bothun G. D., Schombert J. M., 1995, *AJ*, **110**, 573
- Mihos J. C., et al., 2022, *ApJ*, **924**, 87
- Montes M., Trujillo I., Infante-Sainz R., Monelli M., Borlaff A. S., 2021, *ApJ*, **919**, 56
- Montes M., et al., 2024, *A&A*, **681**, A15
- Moreno J., et al., 2022, *Nature Astronomy*, **6**, 496
- Muñoz R. R., Côté P., Santana F. A., Geha M., Simon J. D., Oyarzún G. A., Stetson P. B., Djorgovski S. G., 2018, *ApJ*, **860**, 66
- Müller O., et al., 2020, *A&A*, **640**, A106
- Müller O., et al., 2021, *ApJ*, **923**, 9
- Nandi N., Banerjee A., Narayanan G., 2023, *arXiv e-prints*, p. arXiv:2310.08925
- Navarro J. F., Frenk C. S., White S. D. M., 1997, *ApJ*, **490**, 493
- Olivant T. E., 2015, *Guide to NumPy*, 2nd edn. CreateSpace Independent Publishing Platform, USA
- Oman K. A., Marasco A., Navarro J. F., Frenk C. S., Schaye J., Benítez-Llambay A., 2019, *MNRAS*, **482**, 821
- Pontzen A., Governato F., 2012, *MNRAS*, **421**, 3464
- Prole D. J., et al., 2019, *MNRAS*, **484**, 4865
- Read J. I., Agertz O., Collins M. L. M., 2016a, *MNRAS*, **459**, 2573
- Read J. I., Iorio G., Agertz O., Fraternali F., 2016b, *MNRAS*, **462**, 3628
- Román J., Trujillo I., 2017, *MNRAS*, **468**, 703
- Rong Y., Guo L., Gao L., Liao S., Xie L., Puzia T. H., Sun S., Pan J., 2017, *MNRAS*, **470**, 4231
- Rong Y., Hu H., He M., Du W., Guo Q., Wang H.-Y., Zhang H.-X., Mo H., 2024, *arXiv e-prints*, p. arXiv:2404.00555
- Roper F. A., Oman K. A., Frenk C. S., Benítez-Llambay A., Navarro J. F., Santos-Santos I. M. E., 2023, *MNRAS*, **521**, 1316
- Ruiz-Lara T., et al., 2018, *MNRAS*, **478**, 2034
- Safarzadeh M., Scannapieco E., 2017, *ApJ*, **850**, 99
- Saifollahi T., Zaritsky D., Trujillo I., Peletier R. F., Knapen J. H., Amorisco N., Beasley M. A., Donnerstein R., 2022, *MNRAS*, **518**, 499
- Sales L. V., Navarro J. F., Peñafiel L., Peng E. W., Lim S., Hernquist L., 2020, *MNRAS*, **494**, 1848

- Sandage A., Binggeli B., 1984, *AJ*, **89**, 919
- Sands I. S., et al., 2024, *arXiv e-prints*, p. [arXiv:2404.16247](https://arxiv.org/abs/2404.16247)
- Schombert J. M., Bothun G. D., Schneider S. E., McGaugh S. S., 1992, *AJ*, **103**, 1107
- Shen Z., van Dokkum P., Danieli S., 2021a, *ApJ*, **909**, 179
- Shen Z., et al., 2021b, *ApJL*, **914**, L12
- Shen Z., van Dokkum P., Danieli S., 2023, *ApJ*, **957**, 6
- Sifón C., van der Burg R. F. J., Hoekstra H., Muzzin A., Herbonnet R., 2018, *MNRAS*, **473**, 3747
- Skibba R. A., et al., 2012, *ApJ*, **761**, 42
- Sprayberry D., Impey C. D., Bothun G. D., Irwin M. J., 1995, *AJ*, **109**, 558
- Tollét E., et al., 2016, *MNRAS*, **456**, 3542
- Toloba E., et al., 2018, *ApJL*, **856**, L31
- Toloba E., et al., 2023, *arXiv e-prints*, p. [arXiv:2305.06369](https://arxiv.org/abs/2305.06369)
- Torrealba G., et al., 2019, *MNRAS*, **488**, 2743
- Tremmel M., Wright A. C., Brooks A. M., Munshi F., Nagai D., Quinn T. R., 2020, *MNRAS*, **497**, 2786
- Trujillo I., 2021, *Nature Astronomy*, **5**, 1182
- Trujillo-Gomez S., Kruijssen J. M. D., Keller B. W., Reina-Campos M., 2021, *MNRAS*, **506**, 4841
- Trujillo-Gomez S., Kruijssen J. M. D., Reina-Campos M., 2022, *MNRAS*, **510**, 3356
- Van Nest J. D., Munshi F., Wright A. C., Tremmel M., Brooks A. M., Nagai D., Quinn T., 2022, *ApJ*, **926**, 92
- Venhola A., et al., 2017, *A&A*, **608**, A142
- Venhola A., et al., 2022, *A&A*, **662**, A43
- Villaume A., et al., 2022, *ApJ*, **924**, 32
- Virtanen P., et al., 2020, *Nature Methods*, **17**, 261
- Wang C.-W., Cooper A. P., Bose S., Frenk C. S., Hellwing W. A., 2023, *ApJ*, **958**, 166
- Webb K. A., et al., 2022, *MNRAS*, **516**, 3318
- Willmer C. N. A., 2018, *ApJS*, **236**, 47
- Wolf J., Martinez G. D., Bullock J. S., Kaplinghat M., Geha M., Muñoz R. R., Simon J. D., Avedo F. F., 2010, *MNRAS*, **406**, 1220
- Wright A. C., Tremmel M., Brooks A. M., Munshi F., Nagai D., Sharma R. S., Quinn T. R., 2021, *MNRAS*, **502**, 5370
- Wu Z., Kravtsov A., 2024, *arXiv e-prints*, p. [arXiv:2405.08066](https://arxiv.org/abs/2405.08066)
- Yagi M., Koda J., Komiyama Y., Yamanoi H., 2016a, *ApJS*, **225**, 11
- Yagi M., Koda J., Komiyama Y., Yamanoi H., 2016b, *ApJS*, **225**, 11
- Yozin C., Bekki K., 2015, *MNRAS*, **452**, 937
- Zöllner R., Kluge M., Staiger B., Bender R., 2024, *ApJS*, **271**, 52
- van Dokkum P. G., Abraham R., Merritt A., Zhang J., Geha M., Conroy C., 2015a, *ApJ*, **798**, L45
- van Dokkum P. G., et al., 2015b, *ApJ*, **804**, L26
- van Dokkum P., et al., 2016, *ApJL*, **828**, L6
- van Dokkum P., et al., 2017, *ApJ*, **844**, L11
- van Dokkum P., et al., 2018a, *Nature*, **555**, 629
- van Dokkum P., et al., 2018b, *ApJ*, **856**, L30
- van Dokkum P., Danieli S., Abraham R., Conroy C., Romanowsky A. J., 2019a, *ApJ*, **874**, L5
- van Dokkum P., et al., 2019b, *ApJ*, **880**, 91
- van Dokkum P., et al., 2022a, *Nature*, **605**, 435
- van Dokkum P., et al., 2022b, *ApJ*, **940**, L9
- van der Burg R. F. J., Muzzin A., Hoekstra H., 2016, *A&A*, **590**, A20
- van der Burg R. F. J., et al., 2017, *A&A*, **607**, A79

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